

Atms 502, CSE 566

## Numerical Fluid Dynamics



Description: Simulating the breaking gravity wave and turbulence from the vorticity view of a weak three-dimensional gravity wave bubble and the resulting secondary vortex tubes. The image shows the interaction between the primary gravity wave and the vorticity tubes.


  
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2

ATMS 502  
CSE 566

Tuesday,  
26 March 2019

Class #19

## Plan for Today

- 1) Review
  - Visualization
  - HPC architecture
  - Amdahl's law, parallel performance
- 2) Time differencing
  - including time differencing table
- 3) Nonlinear instability & aliasing
  - problems and methods to mitigate them

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## Time differencing

3

“STAGES” AND “LEVELS”  
HIGHER ORDER ACCURACY

Terminology:

- $n, n+1 \dots$  time levels
- $\alpha, \beta, \dots$  coefficients of spatial derivatives
- $F$  approximations to spatial derivatives

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## Time differencing: stages vs. levels

4

- **Review:** time differencing includes
  - 1) how we express the *time derivative*
  - at what *time levels* we evaluate the *spatial derivatives*
- **Levels**
  - ... refers to how many time levels are in our scheme
  - Lax-Wendroff: *2-level*. Leapfrog: *3-level*
- **Stages**
  - ... refers to how many times we *evaluate the spatial derivatives*
  - Lax-Wendroff: *single-stage*. Runge-Kutta: *2 or more stages*

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5

**Summary:** Single-stage, 2-level schemes

- **Single-stage:** evaluate spatial derivatives once
- **2-level:** there are two time levels,  $n$  and  $n+1$

$$\frac{\phi^{n+1} - \phi^n}{\Delta t} = \alpha F(\phi^n) + \beta F(\phi^{n+1}), \quad \alpha + \beta = 1$$

**AMPLITUDE**

Trapezoidal best (T overlaps E)

**PHASE**

Trapezoidal best (1/4 error of others)

Euler method:  $\alpha=1, \beta=0$   
 Backward method:  $\alpha=0, \beta=1$   
 Trapezoidal method:  $\alpha=\beta=1/2$

$$|A| = 1 - (\alpha - \beta)^2 \left( \frac{\omega^2 \Delta t^2}{1 + \beta \omega^2 \Delta t^2} \right)$$

Schemes: Forward F • Backward B • Trapezoidal T • 2<sup>nd</sup>-order Runge-Kutta R • Matsuno M

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6

**Summary:** Single-stage, 3-level schemes

- **Single-stage:** evaluate spatial derivatives once
- **3-level:** there are 3 time levels,  $n-1, n, n+1$

$$\frac{\phi^{n+1} - (\alpha_1 \phi^n + \alpha_2 \phi^{n-1})}{\Delta t} = \beta_1 F(\phi^n) + \beta_2 F(\phi^{n-1})$$

$\alpha_1 = 1 - \alpha; \beta_1 = \frac{\alpha_2 + 3}{2}, \beta_2 = \frac{\alpha_2 - 1}{2}$   
 ( $\beta$  restrictions make schemes at least 2<sup>nd</sup>-order – Durran p. 58)

- **Leapfrog:**  $\alpha_1=0, \alpha_2=1, \beta_1=2, \beta_2=0$  Time filtering: comp. mode; even/odd ..
- **Leapfrog-trapezoidal** A predictor-corrector method.  
 $\phi^n = \phi^{n-1} + 2\Delta t F(\phi^n)$   
 2<sup>nd</sup>-order unlike time-filtered LF; computational mode damped
- **Adams-Bashforth**  
 $\phi^{n+1} = \phi^n + (\Delta t) \left[ \frac{3}{2} F(\phi^n) - \frac{1}{2} F(\phi^{n-1}) \right]$   
 2<sup>nd</sup>-order, fwd time, weak instability; computational mode damped. Higher-order versions of A-B exist.

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7

**Overview:** Multi-stage (RK)/step methods

- **Objective:** higher-order accuracy in time
- **Multistage:** Durran § 2.3: evaluate spatial  $F$  terms at several times between  $n\Delta t, (n+1)\Delta t$
- **Multistep:** Durran § 2.4: information from prior levels incorporated in integration formula.
  - Multistep: extra storage needed, but fewer evaluations of  $F$
  - Multistage or multistep: computational modes arise.
  - Multiple forms of RK (Runge-Kutta) methods exist, e.g. low-storage form.
- **More steps/or/stages: more (computational) work.**
  - Good: less restrictive time step. Bad: more computational modes.
  - For more information: see Durran § 2.3.2, RK3 and RK4 methods.

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8

**Overview:** Higher time accuracy in 3-D

- **Problem:** order of directional operators matters.
- **Theory:** Durran § 4.3: compares exact (that from Taylor series) vs. finite difference results
  - Let's call X-operator  $F_1, \dots$  and ... Y-operator  $F_2, \dots$
  - Then if  $F_1 F_2 = F_2 F_1$ , we say the operators *commute*
- **Plan:** if operators *don't* commute –
  - We can still get higher temporal accuracy ...
  - Use *Strang Splitting*. In 3D: Durran eq. 4.59 p. 172 –

$$\left[ F_1(\Delta t/2) \right] \left[ F_2(\Delta t/2) \right] \left[ F_3(\Delta t) \right] \left[ F_2(\Delta t/2) \right] \left[ F_1(\Delta t/2) \right]$$

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## Time differencing summary: Durran § 2.6

⑨

- Durran table 2.1-2.2: summary of methods
- Order of scheme
- Storage factor
  - ✦ Number of full arrays needed
  - ✦ Not given if implicit; depends on method
- Efficiency factor
  - ✦ Largest stable step  $\Delta t$  by # evaluations of  $\mathbf{F}$
- Max s
  - ✦ largest stable  $\omega\Delta t$ !!

Table 2.1 Summary of methods for the solution of ordinary differential equations. The second- and third-order Runge-Kutta methods are two-stage methods.  $k = \Delta t$  Durrant 2<sup>nd</sup> ed. TABLES 2.1-2.2, pp. B3-B4

Method	Order	Formula	Storage factor	Amplification factor	Phase error	Max s
Forward	1	$\phi_{n+1} = \phi_n + hF(\phi_n)$	2	0	$1 + \frac{h^2}{2}$	1 - $\frac{h^2}{3}$
Backward	1	$\phi_{n+1} = \phi_n + hF(\phi_{n+1})$	2	0	$1 + \frac{h^2}{2}$	1 - $\frac{h^2}{3}$

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## Aliasing, nonlinear instability, and conservation

⑩

REFERENCES:  
 DURRAN SECTION § 4.4.1, 4.5  
 ROBERT WILHELMSON NOTES  
 HALTNER AND WILLIAMS SECTION 5-1-1-1  
 PAUL SCHOPF NOTES, SCHOOL OF COMPUTATIONAL SCIENCES,  
 GEORGE MASON UNIVERSITY (MASON.GMU.EDU)

## (Nonlinear) advection

⑪

- Back to the familiar.

$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} = 0 \quad \text{or} \quad \frac{\partial}{\partial t} \left( u^2 \right) = 0$$

○ The problem for finite differences lies in the multiplication.\*

- Let  $u = \sin(kx)$ ; then

$$u = \sin kx \Rightarrow \frac{\partial}{\partial x} (u^2) = k \sin 2kx$$

○ What effective wavenumber are we working with now?

## Aliasing

⑫

- Even if  $u = \sin kx$  is resolved...

- If ...  $(k_{max}/2) < k < k_{max}$ ,
- the nonlinear term puts energy into:

$$k_{max} < k < 2k_{max}$$

- This is not resolvable!

- ✓ It will appear in other, resolvable wavelengths.
- ✓ This is aliasing.

### Aliasing: Example

13

- Some studies have investigated **gravity waves** near weather fronts.
- They appeared in observations
- They appeared in **high-resolution** simulations.
- Some of the *modeled* waves were not real.
- They appeared at **high horizontal resolution**.

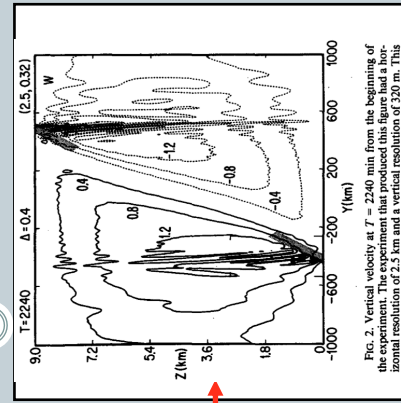


FIG. 2. Vertical velocity at T = 2240 min from the beginning of the experiment. The experiment that produced this figure had a horizontal resolution of 2.5 km and a vertical resolution of 330 m. This

GALL ET AL. 1988

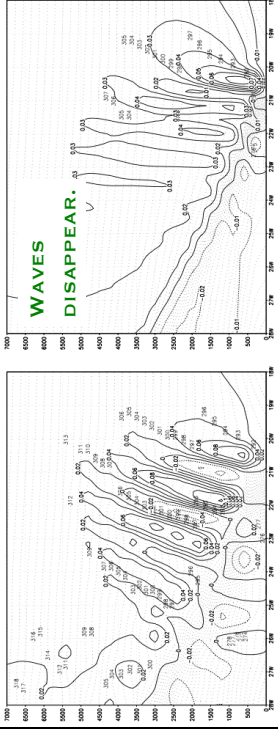
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### Spurious waves

15

GLEVEL=9, Δz=600m



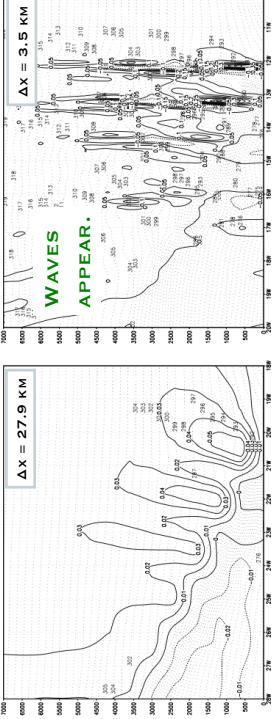
- Iga (2005) - waves on a front. Vertical velocity is shown.  $\Delta x=14$  km. Only change is decrease in  $\Delta z$ .

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### Spurious waves

14



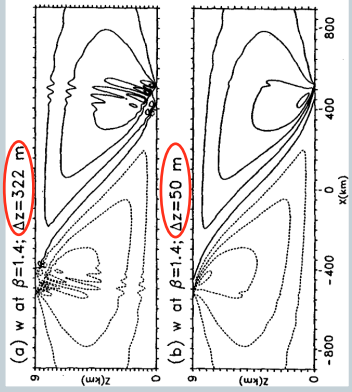
- Iga (2005) - waves on a front. Vertical velocity is shown.
- Only change is decrease in  $\Delta x$ .

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### Spurious waves

16



- Snyder et al. 1993. Only change is  $\Delta z$ .
- Waves appear at high **horizontal** resolution
- Waves **disappear** at increased **vertical** resolution.
- Why?**

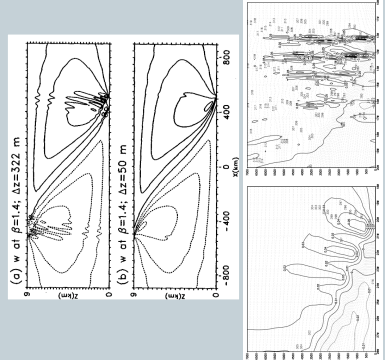
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### Spurious waves

17

- One possibility ...
- Reducing  $\Delta x$  while keeping  $\Delta z$  constant
- results in small horizontal waves with small vertical scales ... too small to be resolved.
- Uh-oh!

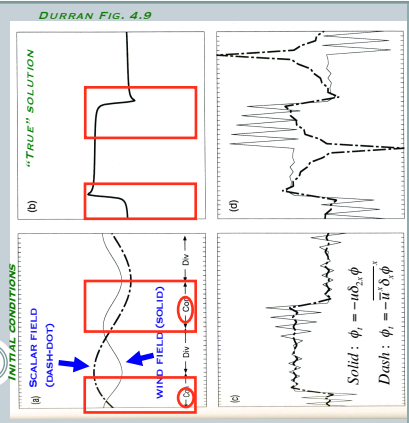


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### Aliasing

18

- Simple differential (t) - difference (x) equation results
- Problem: linear equation, variable coefficients in other words,  $c=c(x)$ , but not (t)
- Small scales grow preferentially



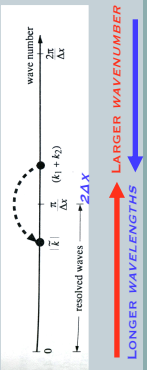
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### Aliasing

19

- Where do unresolved waves "go" ?
- Durrnan eqn. 3.91:
- Examples:
  - $\times 2\Delta x \rightarrow 2.5\Delta x \rightarrow 10\Delta x$
  - $\times (4/3)\Delta x \rightarrow 4\Delta x$
- Note if both waves are  $4\Delta x$  or longer: **no aliasing**

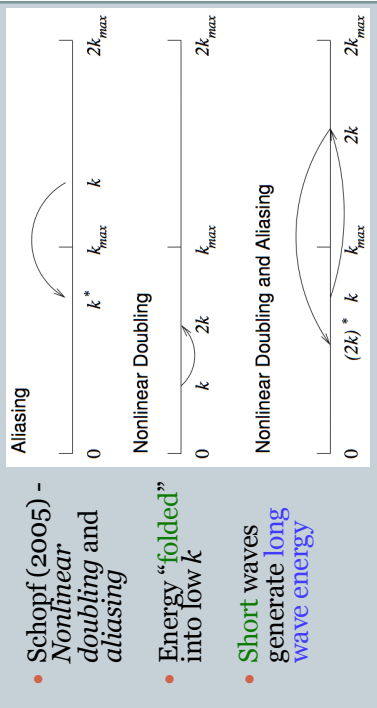
$$\tilde{k} = \begin{cases} k_1 + k_2 - \frac{2\pi}{\Delta x}, & \text{if } k_1 + k_2 > \frac{\pi}{\Delta x} \\ k_1 + k_2 + \frac{2\pi}{\Delta x}, & \text{if } k_1 + k_2 < \frac{-\pi}{\Delta x} \end{cases}$$



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### Aliasing

20



- Schopf (2005) - Nonlinear doubling and aliasing
- Energy "folded" into low k
- Short waves generate long wave energy

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### Nonlinear instability

(21)

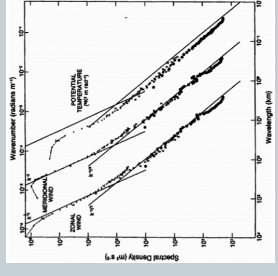
- If energy flows into wavenumbers *just above*  $k_{max}$  ...
- Aliasing folds it back to wavenumbers *just below*  $k_{max}$ 
  - builds up energy near grid resolution limit
  - Further nonlinear interaction enhances flow into wavenumbers just above  $k_{max}$ : accelerates process
  - This is nonlinear instability.
- What about amplitude?
  - More energy at low  $k$ . Say  $k_2$  just below  $k_{max}$  ...
  - Nonlinear ( $k_1+k_2$ ) has more energy *if*  $k_1$  *small*.



### Nonlinear instability

(22)

- $2\Delta x \cdot 2.5\Delta x \rightarrow 10\Delta x$
- $2\Delta x \cdot 8\Delta x \rightarrow 2.7\Delta x$
- more energy: *small*  $k$



### Aliasing in wavenumber space

(23)

Input wave	Input $k / k_{max}$	+ $2\Delta x$ $k / k_{max}$	aliased $k / k_{max}$	= final $N\Delta x$
2.1 $\Delta x$	0.95	1.95	0.05	42 $\Delta x$
5.0 $\Delta x$	0.40	1.40	0.60	3.3 $\Delta x$
10 $\Delta x$	0.20	1.20	0.80	2.5 $\Delta x$
20 $\Delta x$	0.10	1.10	0.90	2.2 $\Delta x$

### Aliasing in wavenumber space

(24)

