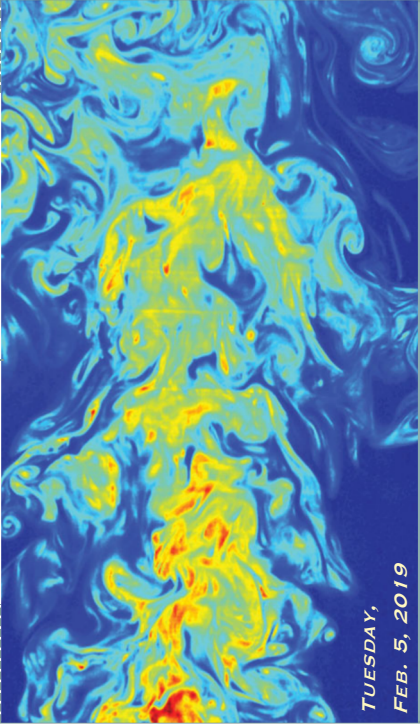


Atms 502, CSE 566

Numerical Fluid Dynamics



**TUESDAY,
FEB. 5, 2019**

*ATMS 502 - Spring 2019 Prerequisite of the turbulent/non-turbulent interface of a jet - Westermann, Fukushima, Pedersen and Hama 2019

Plan for Today

- 1) REVIEW Takacs vs. Polar plots
- 2) NUMERICAL METHODS :
 - Stability, continued:
 - ✦ Norms; von Neumann's method
 - ✦ Apply to a numerical method
 - ✦ Operator definitions
 - Phase error
 - ✦ *wavelength-dependent phase speeds*
 - The modified equation
- 3) CODE/DATA: Program #2 - continued

*Homework #1
prob #4 typo:
should be
2Δx not Δx !!*

ATMS 502
CSE 566

Tuesday,
5 February 2019
Class #7

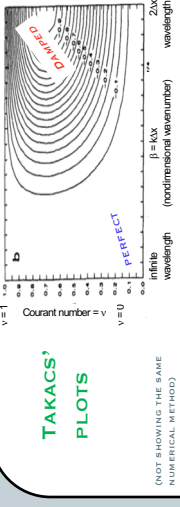
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Review: Plots of scheme behavior

AMPLITUDE

TAKACS' PLOTS



Countant number = v

$\beta = k\Delta x$
(non-dimensional wavenumber)

infinite wavelength $2\Delta x$ wavelength

$v = 1.25$ **AMPLIFYING**

$v = 1.0$ **PERFECT**

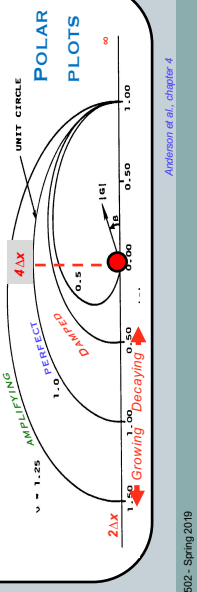
$v = 0.5$ **DAMPED**

$v = 0.25$ **Growing**

$v = 0.125$ **Decaying**

UNIT CIRCLE

POLAR PLOTS



$|e|$

1.0 **Growing**

0.5 **Decaying**

0.25

0.125

1.0 **PERFECT**

0.5 **DAMPED**

0.25 **AMPLIFYING**

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Anderson et al., Chapter 4

Stability

4

OBJECTIVES:

DEVELOP THEORY AND METHODOLOGY FOR DETERMINING IF, HOW, AND WHEN A SCHEME HAS SATISFACTORY STABILITY.

FOLLOWING NOTES HANDED OUT IN LAST CLASS !!

References:

- A009 – Instability (physical)
- C015 – Instability (numerical)

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Operator definitions

USED FOR REMAINDER OF CLASS

FOLLOWING NOTES HANDED OUT IN LAST CLASS !!

ATMS 502 - Spring 2019 C032 - Operator notation for finite differences 2/5/19

Dispersion & phase error

IN CONTEXT OF: FOURIER SERIES

References:

- C016 – Fourier series
- C023 – Dispersion
- C033 – Gibbs phenomenon

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Fourier series

- Assuming we have a *periodic* function of period 2π , we can represent it with the following trig series:
- Solve using the *Euler Formulas*:

$$f(x) = a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx)$$

Period=T

$$a_0 = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) dx$$

$$a_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \cos(nx) dx$$

$$b_n = \frac{1}{\pi} \int_{-\pi}^{\pi} f(x) \sin(nx) dx$$

$$a_0 = \frac{1}{T} \int_{-T/2}^{T/2} f(t) dt$$

$$a_n = \frac{2}{T} \int_{-T/2}^{T/2} f(t) \cos \frac{2n\pi t}{T} dt$$

$$b_n = \frac{2}{T} \int_{-T/2}^{T/2} f(t) \sin \frac{2n\pi t}{T} dt$$

$$n = 1, 2, \dots$$

ATMS 502 - Spring 2019 C016: Fourier Series 2/5/19

Fourier series

- Let's look at the representation of a square wave.
- The analytical solution turns out to be:

$$f(x) = \begin{cases} -k & \text{when } -\pi < x < 0 \\ k & \text{when } 0 < x < \pi \end{cases}$$

$$a_0 = 0$$

$$a_n = 0$$

$$b_n = \frac{4k}{n\pi}$$

$$n = 1, 3, 5, \dots \text{ such that } f(x) \approx \frac{4k}{\pi} \left(\sin x + \frac{1}{3} \sin 3x + \frac{1}{5} \sin 5x + \dots \right)$$

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Fourier series & square wave

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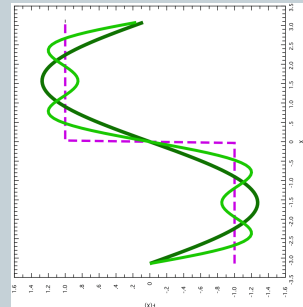
- Fourier series *cannot be used (with the strict equality) for a discontinuity; doing so invokes the **Gibbs phenomenon**.*
- “While the rms error of the Fourier series goes to zero for an infinite number of terms, equality at every point is not guaranteed; the **Gibbs phenomenon peaks have finite height and zero width.**”
- Coefficients in a truncated Fourier series can be used to reduce or eliminate the Gibbs effect; see References for more on these windowing methods.

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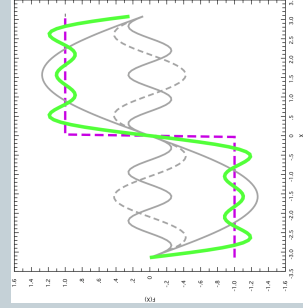
C023 - Gibbs phenomenon

2/5/19

- Using one wave vs. sum of two waves



- Sum of first three waves. A better approximation.



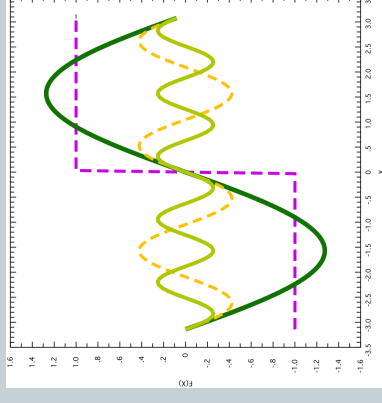
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- $N=1,3,5$: **three waves shown:** $\sin(x)$, $\sin(3x)$, $\sin(5x)$...



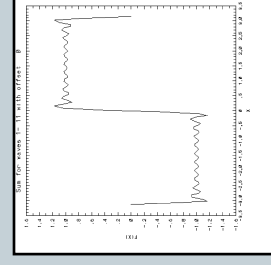
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What if our scheme has phase errors?

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- We'll take **one component - 1 harmonic** - and change it
- The second wave -- with **wave number $k=3$** -- will have a **phase shift** added to it.
- There are 11 contributing waves; all other components - the other 10 waves - will remain the **same**.



In this example, there is no phase shift - we add the terms correctly.

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C023 - Dispersion - numerical phase error

2/5/19

Phase errors

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- Dispersion distorts the solution by introducing **wavenumber-dependent phase errors**
- The **phase relationship** between the different wave components is changed.
- In this example, the 2nd wave component was shifted by +45°

Sum for waves 1-11 with offset 45

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The modified equation

(14)

WHY WE SEE CHARACTERISTIC (OR REALLY, DOMINANT) ERRORS FOR TYPES OF SCHEMES

ATMS 502 - Spring 2019 C034: Modified equation 25/19

The modified equation

(15)

- Upstream scheme - truncation error:

$$u_t + cu_x = -\frac{\Delta t}{2} u_{tt} + \frac{c\Delta x}{2} u_{xx} + \text{higher order terms} \quad (1)$$
- To say more about this scheme, we want the right side to be in terms of **X derivatives, only.**
 - Take equation (1), and take d/dt:

$$u_{tt} + cu_{xt} = -\frac{\Delta t}{2} u_{ttt} + \frac{c\Delta x}{2} u_{xtx} + \text{higher order terms} \quad (2)$$
 - And now take -c*d/dx of (1):

$$-cu_{xt} - c^2 u_{xx} = \frac{c\Delta t}{2} u_{xt} - \frac{c^2 \Delta x}{2} u_{xxx} + \text{higher order terms} \quad (3)$$

ATMS 502: Spring 2019 • C022: Amplitude error • C023: Phase error • C034: Modified equation 25/19

Modified equation: implicit viscosity

(16)

- Our original equation:

$$u_t + cu_x = -\frac{\Delta t}{2} u_{tt} + \frac{c\Delta x}{2} u_{xx} + \text{higher order terms} \quad (1)$$
- Our expression for u_{tt} was:

$$u_{tt} = c^2 u_{xx} + \Delta t \left(-\frac{u_{ttt}}{2} + \frac{c}{2} u_{xtt} + \dots \right) + \Delta x \left(\frac{c}{2} u_{xt} - \frac{c^2}{2} u_{xxx} + \dots \right) \quad (4 = \text{eq. 2}) + (3)$$
- Substituting, we get:

$$u_t + cu_x = \frac{c\Delta x}{2} (1 - \nu) u_{xx} + (\dots)(\Delta x)^2 u_{xxx} + \text{higher order terms}$$
 - This is the **modified equation**
 - It is what is **actually solved** by the F.D. method
 - U_{xx} term: **"implicit" artificial viscosity!** (not an implicit numerical method)

ATMS 502 - Spring 2019 C037: Implicit viscosity 25/19

Modified equation: dissipative errors

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- Modified equation:

$$u_t + cu_x = \frac{c\Delta x}{2}(1 - \nu)u_{xx} + \text{higher order terms}$$

- Why is this diffusive?
 - Recap: we started with a **hyperbolic** PDE, $u_t + cu_x = 0$.
 - The modified equation tells us what we **really** solving.
 - Now our analysis reveals terms like: $u_t = 0 + \nu u_{xx}$
 - This is a **parabolic** equation!
 - wait, weren't we solving a **hyperbolic** (transport) equation?
 - Parabolic > for example, heat transfer
 - Bottom line: Dissipation!

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Diffusive upstream method: polar plot

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- Modified equation for the upstream method:

$$u_t + cu_x = \frac{c\Delta x}{2}(1 - \nu)u_{xx} + \text{higher order terms}$$

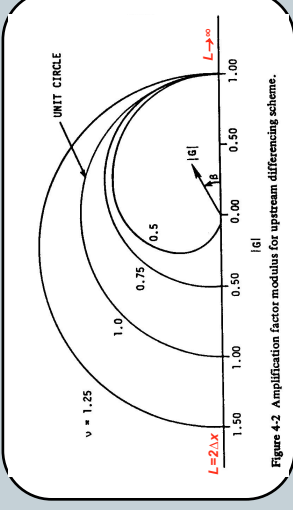


Figure 4.2 Amplification factor modulus for upstream differencing scheme.

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Shift condition

(19)

- Modified equation for the upstream method:

$$u_t + cu_x = \frac{c\Delta x}{2}(1 - \nu)u_{xx} + \text{higher order terms}$$

- What if the Courant number $\nu=1$ here?
 - The u_{xx} term on the right side of the modified equation disappears
 - In fact, we know this scheme has a **shift condition** when the Courant number $\nu=1$.
 - this means the solution is shifted one grid point per time step.

$$u_j^{n+1} = u_j^n - \nu(u_j^n - u_{j-1}^n) ; \nu = \left(\frac{c\Delta t}{\Delta x} \right)$$

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C035: Shift condition

2/5/19

Problems with upwinding-type schemes

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- Note the problem with upstream-type methods:
 - The implicit diffusion u_{xx} from this scheme is **dependent on the local flow speed**.
 - So you have uneven damping throughout your flow
 - Some modelers rely on this **implicit*** damping to stabilize their solution. We'll return to this later.

* Again, *implicit* here refers to the damping, not an implicit numerical method.

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C036: Upwind advection schemes; C037: Implicit viscosity

2/5/19

Summary: Errors vs. order of accuracy



- **Explicit artificial viscosity**
 - An added term in a difference equation *designed* to add damping
- **Implicit artificial viscosity**
 - Unphysical damping as a *consequence* of the finite difference scheme
- **Summary of error properties**
 - Dispersion - **distortion of waves** as a result of **odd** derivative terms in the truncation error - **even** order accuracy!
 - Dissipation - **reduction of gradients** as a result of **even** derivative terms in the truncation error - **odd** order accuracy!
 - Diffusion: combined effect of dissipation and dispersion

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C037: Implicit viscosity

2/5/19

Convergence

(22)

- We've discussed:**
- Finite difference approximations to derivatives
 - Truncation error
 - Taylor series
 - Order of accuracy
 - Consistency
 - Stability
 - Von Neumann's method
 - CFL
- Convergence -**
Lax equivalence theorem
 If a finite difference scheme is:
- Linear
 - Stable
 - Accurate of order $(\Delta t)^p$, $(\Delta x)^q$, then
- It is **convergent** of order (p,q)
- This has not been shown to apply to nonlinear PDE's !!
 - Durran, §2.1.3, p. 40

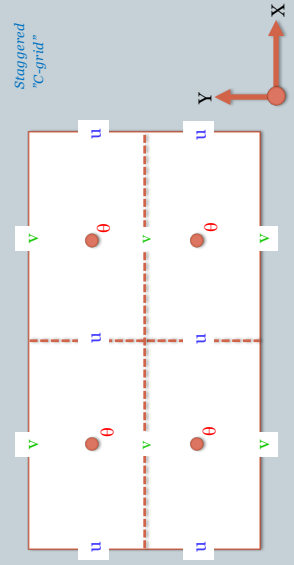
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C011: Convergence

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Computer Program 2

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C037: Implicit viscosity

2/5/19

Program 2: Advection

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- **Advection**
 - I set up 1-D arrays in my advection routines -
 - × $q1d(0:nx+1)$, $u1d(nx+1)$, $v1d(ny+1)$ no ghost points for U, V !!
- Advecting rows (X)
 - copy $q1(i,j)$ to $q1d$
 - copy $u(i,j)$ to $u1d$
 - pass $q1d$, $u1d$ to **advect1d**
 - × advect1d returns $q1d_out$
 - copy $q1d_out$ to $q1(i,j)$
- Advecting columns (Y)
 - copy $q1(i,j)$ to $q1d$
 - copy $v(i,j)$ to $v1d$
 - pass $q1d$, $v1d$ to **advect1d**
 - × advect1d returns $q1d_out$
 - copy $q1d_out$ to $q1(i,j)$

discuss: how many 2D q() arrays here?

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C014: Directional splitting

2/5/19